# Construction of velocity similarity graph for velocity field analysis in granular intruder motion

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#### Abstract

We present a method based on graph community detection algorithms to analyse velocity fields induced by an intruder particle impinging upon a stationary bed of particles. Based on velocity relations between the pairs of adjacent particles, the "velocity similarity" graphs are built where the graph vertices represent the particles and the edge weights are calculated according to the velocities of the respective particle pairs. A few different expressions for the edge weights are tested. Based on the graph, a Louvain community detection algorithm with the "geographic" null model is used to identify the goups of particles moving in a coordinated manner, represented in the graph as a community of vertices, for which the community detection algorithms developed for graph analysis can be applied. Selection of the expression of the graph edge weights based on the velocities of the respective particles influences the resulting graph structure and thereby has an influence on the community detection results.

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# **1** INTRODUCTION

Mechanical behaviour of granular media is determined by multiple interactions of the constituent particles giving rise to the bulk properties in moving matter. Discrete element simulations (DEM) provide the data of each individual particle i at certain time moments, i.e., the "microscopic" state of the matter, whereas the bulk behaviour, encountered in experiments and practical applications involving granular materials, depends on the large groups of particles arranged in certain structural patterns that can be thought of as a "macroscopic" state. Identification of characteristic larger-scale features of motion of granular media from the particle data is important for understanding the bulk behaviour.

The processes in granular media are based on pairwise particle interactions. These interactions can be represented as a graph where the vertices represent the particles and their interactions as the edges. Earlier, graph theory algorithms were applied for analysis of larger scale structures in granular media [1], as the force chains can be readily represented by graphs. Application of certain algorithms from graph analysis, such as community detection [2, 3], becomes equivalent to identification of structures in granular media [4, 5, 6]. Analogously, groups of particles moving in a coordinated manner can be identified as graph communities if the edge weights represent the velocity relations [7].

Motion of an "intruder" particle in granular media was analysed from a number of viewpoints [8, 9, 10, 11, 12, 13, 14, 15, 16]. In our case, it is interesting because it provides a characteristic feature of a "shock wave" with rather clear boundary propagating in front of the intruder particle. Detection of this wave can be used as a test of the proposed method. The velocity similarity graph can be constructed using various expressions of the edge weight. We analysed a few expressions to calculate the graph edge weights to construct such a graph from the velocities of the respective particles.

# 2 SIMULATION METHOD

The considered system consists of a packed bed of 10000 spherical particles partially filling a container (Fig. 1a). The bottom corners of the container are slanted, in order to facilitate the particle motion and thereby prevent excessive mechanical force buildup, in order to avoid possible instabilities during the DEM simulations. The particle radii are random in the range from  $r_{\rm min} = 0.005$  m to  $r_{\rm max} = 0.015$  m. The particle density is  $\rho = 2500$  kg/m<sup>3</sup>. Additionally, an extra particle ("intruder") with the radius  $r_{\rm int} = 0.1$  m and density  $\rho_{\rm int} = 5000$  kg/m<sup>3</sup> is placed at a small height above the packed bed. The initial "intruder" velocity is 10 m/s directed vertically downward. As the intruder reaches the packed bed, the particles therein are put to motion by contact with the intruder (Fig. 1). For easy visualisation, all the particles are fixed in a single plane, but their shapes are considered spherical (a quasi-2D problem). The process was simulated using the ligghts software package [17, 18]. In the initial stages, the particles nearest to the intruder start to move and, as the intruder continues its passage, the motion gradually propagates throughout the packed bed and a boundary between the moving and immobile particles appears that can be thought of as a kind of shock wave propagation; even though it is not a rigorous definition, it is a convenient term to refer to the areas of moving particles as opposed to the mostly immobile bulk. Later, the particles in the whole volume of the packed bed start to move and the character of the motion changes. This early propagation of the "shock wave" is a characteristic feature in the granular bed, containing a rather clearcut single particle group, that can be thought of as a group of particles sharing a certain characteristic parameter (in our case – speed). At later stages of the intruder motion, a canal opens in the wake of the intruder particle (Fig. 1d) that gradually backfills later on (Fig. 1e,f). In these later stages, the character of particle motion is more complex, encompassing a few areas where the nature of motion (direction and speed) varies.

The emergence of particle groups moving faster than the surrounding bulk, i.e., the structure of the velocity field, can be estimated by visualising only the particles moving faster than a certain predefined threshold. Fig. 2 shows the velocity field (particle velocity vectors) for particles moving faster than a certain fraction of the maximum velocity  $|\mathbf{v}|_{\max} \approx 11 \text{ m/s}$ , at the time moment t = 0.7 s. A group of particles having velocities  $> 0.05 |\mathbf{v}|_{\max}$  is visible in front of the advancing intruder particle. The velocity field configuration at this time moment is convenient for analysis, because the single group of moving particles, in contrast to mostly immobile particles located deeper, can be used as a benchmark for easy estimation of detection methods of such groups. At this time, the group still has not reached the walls of the container, therefore, there is no "secondary" motion due to particles bouncing off the walls, and the velocity field has a rather regular structure.

At a later time, after the intruder particle has reached the bottom of the container, the "ballistic canal" starts to fill in. By the same thresholding, a few groups of particles moving bilaterally into the void of the canal can be seen at t = 3, 5 s (Fig. 3).

# 3 BUILDING THE VELOCITY RELATION-SHIP GRAPHS

As we represent the similarity of velocities of the nearby particles as the graph edges, the groups of particles having similar velocities can be detected using the mentioned method based on the "velocity similarity" graph.

The graph was built based on the velocities  $\mathbf{v}_i$ ,  $\mathbf{v}_j$  of the nearby particles i, j. The particles are considered nearby if  $|\mathbf{x}_i - \mathbf{x}_j| \leq f_R (r_i + r_j)$ , where  $\mathbf{x}_i$  and  $r_i$  are position and radius, respectively, of particle i. To allow inclusion of particle pairs that are not in direct contact but still close enough from each other, an "extension factor"  $f_R = 1.2$  was introduced. The building of the velocity relationship graph is based on defining the graph edge weights depending on the velocities of the respective particles. The particles are considered having "similar" velocity if the angle between their respective velocity vectors is small and the vector lengths are similar. In this case, the values of the scalar (dot) product of the respective vector pairs or the cosine of the angles between these vectors would be high. Moreover, taking into account the speeds of the particles  $|\mathbf{v}|$ , would enhance the groups of swiftly moving particles against the background of mostly immobile bulk. With these considerations in minds, we used the following expressions for the edge weights:

• similarity of velocities:

$$w_{ij}^{\text{sim}} = 1 - \frac{|\mathbf{v}_i - \mathbf{v}_j|}{\max_{(i,j) \in [1, N_p]} |\mathbf{v}_i - \mathbf{v}_j|},\tag{1}$$

• velocity difference:

$$w_{ij}^{\text{diff}} = \frac{|\mathbf{v}_i - \mathbf{v}_j|}{\max_{(i,j)\in[1,N_p]} |\mathbf{v}_i - \mathbf{v}_j|};\tag{2}$$



Figure 1: Passage of the intruder at different time moments. The particles are colored according to their speeds (velocity moduli  $|\mathbf{v}_i|$ ).



Figure 2: Particle velocities at different speed thresholds at t = 0.7 s: only the velocity vectors for particles with speeds larger than the threshold values  $|\mathbf{v}|_{\rm th}$  are shown, with the threshold values of 0.01 (a), 0.05 (b), 0.10 (c) and 0.15 (d) of the maximum speed  $|\mathbf{v}|_{\rm max} \approx 11$  m/s.



Figure 3: Particle velocities at different speed thresholds at t = 3, 5 s: only the velocity vectors for particles with speeds larger than the threshold values  $|\mathbf{v}|_{\rm th}$  are shown, with the threshold values of 0.01 (a), 0.05 (b), 0.10 (c) and 0.15 (d) of the maximum speed  $|\mathbf{v}|_{\rm max} \approx 11$  m/s.

in the latter case, the edge weights are larger for the particle pairs with differing velocities. This expression can be used to detect the boundaries between the areas with different character of motion, e.g., the shear zones between the streams with different velocities;

• velocity dot product:

$$w_{ij}^{\text{dot}} = \mathbf{v}_i \cdot \mathbf{v}_j,\tag{3}$$

• normalized velocity dot product:

$$w_{ij}^{\text{dot},v} = \frac{\mathbf{v}_i \cdot \mathbf{v}_j}{\max\left(|\mathbf{v}_i|, |\mathbf{v}_j|\right)},\tag{4}$$

• cosine of the angle  $\phi_{ij}$  between the velocity vectors  $\mathbf{v}_i, \mathbf{v}_j$ :

$$w_{ij}^{\cos} = \frac{1}{2} \left( 1 + \cos \phi_{ij} \right), \tag{5}$$

• normalised cosine of the angle between the velocity vectors:

$$w_{ij}^{\cos,v} = \frac{1}{2} \left( 1 + \cos \phi_{ij} \right) \cdot \frac{1}{2} \left( |\mathbf{v}_i| + |\mathbf{v}_j| \right).$$
(6)

The expression (5) is based on the assumption that the particles moving in a coordinated manner will move in similar directions, therefore, the cosine of the angle between their velocity vectors would have higher values than that for particles moving in different directions. For convenience of calculations, the expression (5) is normalised to the range [0, 1]. This definition does not take into account the speeds of the particle pairs, therefore, it would not discern slower and faster particle groups. To enhance detection of faster particle groups, the expression (6) includes multiplication by the average of the respective particle speeds. Similarly, the pairs of velocity vectors of particles moving in similar directions would produce higher values of the scalar product. On the other hand, dot product of velocities of faster particles moving at more different directions would still produce a larger value of the dot product than that of slow particles moving in less different directions. This effect can be avoided by normalisation to the particle speed (Eq. 4).

As seen above, the time at around t = 0.7 s is convenient for analysis because a single rather well- defined structure appears in the velocity field that allows also for visual estimation of the graph structure. The resulting graphs constructed using the expressions (1)–(6) for the time moment t = 0.7 s are shown in Fig. 4a–f. In all the graphs, the edge weights  $w_{ij}$  were normalised to fit to the range  $0 \le w_{ij} \le 1$ .

The depictions of graphs shown in Fig. 4 are not very revealing visually due to abundance of graph edges, most of which have rather small weights in the area where the particles are not yet put to motion. It is therefore more illustrative to show only the edges having weights exceeding a certain threshold. Fig. 5 shows the edges of the graph, built from the particle velocities at t = 0.7 s

using the expression (6), exceeding the values of  $0.05w_{\text{max}}$  (a) and  $0.15w_{\text{max}}$  (b), where  $w_{\text{max}}$  is the maximum edge weight value for this graph. At the low value of threshold  $0.05w_{\text{max}}$ , a localised structure is already visible that can be estimated as similar to the velocity wave seen in the velocity field (Fig. 2). At higher values of the threshold, this structure is gradually eroded; however, most of the low-weight graph edges corresponding to largely immobile particles can be filtered out.

In comparison, the edge weight expression (Eq. 5) leads to a rather uniform edge weight distribution (Fig. 6). Even at a larger threshold of  $0.6w_{\text{max}}$ , distribution of the included edges (with  $w_{ij} \ge 0.6w_{\text{max}}$ ) does not reveal any localised structures. This result can be expected, because only the angles between the motion (even slow) directions of nearby particles are taken into account, therefore, the fast moving particles do not stand out in the overall distribution.

On the contrary, the edge weight calculation based on the dot product of the velocity vectors allows for domination of the fastest particle groups (Fig. 7). In comparison to the case of graph edge weights defined using Eq. 6, a rather small particle group emerges. Normalisation to the maximum velocities of the respective particle pairs (Eq. 4) decreases this contrast and the edges distributed throughout the bulk of the particle bed remain present even at a larger threshold of  $0.6w_{\rm max}$  (Fig. 8).

## 4 RESULTS AND DISCUSSION

Having built the appropriate graph, the standard community detection algorithms known from the graph analysis can be applied [2, 3]. A notable difference between the general community detection approach and that applicable to the graphs built specifically for analysis of granular systems is the choice of null-model: the general algorithms assume that every vertex in the graph has an equal probability to be connected to any other vertex in the null-model. In the granular interaction graphs, where the vertices represent granular particles, this assumption does not hold: only the nearest particles can be in contact, and this constraint should hold also for the representative graph. For this purpose, a "geographic" null model was proposed accounting for this constraint [19, 6]. We used the Louvain algorithm [20] for community detection. The implementation was based on publicly available MATLAB script [21], but we modified it to use the "geographic" null model and to work with GNU Octave. The exhaustive analysis of all the edge weight expressions for building the velocity graph and the resulting community detection results is outside the scope of the present contribution; rather, we focused on the edge weight expression defined by Eq. 6, because it was expected to give the best results based on the above considerations, and this was confirmed by the initial trials.

The community detected in this velocity graph at the time moment t = 0.7 s is shown in Fig. 9 as a green outline; the contour of this particle group is obtained by connecting the centers of the outermost particles in the group. For comparison, the outline is superimposed on the velocity fields shown at the



Figure 4: Edges of graphs of velocity relations during the intruder passage at t = 0.7 s using the different expressions for graph edge weights: a)  $w_{ij}^{\text{sim}} = 1 - \frac{|\mathbf{v}_i - \mathbf{v}_j|}{\max_{(i,j)\in[1,N_p]}|\mathbf{v}_i - \mathbf{v}_j|}$  b)  $w_{ij}^{\text{diff}} = \frac{|\mathbf{v}_i - \mathbf{v}_j|}{\max_{(i,j)\in[1,N_p]}|\mathbf{v}_i - \mathbf{v}_j|}$ , c)  $w_{ij}^{\text{dot}} = \mathbf{v}_i \cdot \mathbf{v}_j$ , d)  $w_{ij}^{\text{dot},v} = \frac{\mathbf{v}_i \cdot \mathbf{v}_j}{\max(|\mathbf{v}_i|, |\mathbf{v}_j|)}$ , e)  $w_{ij}^{\text{cos}} = \frac{1}{2}(1 + \cos\phi)$ , f)  $w_{ij}^{\cos,v} = \frac{1}{2}(1 + \cos\phi) \cdot \frac{1}{2}(|\mathbf{v}_i| + |\mathbf{v}_j|)$ . The graph edges are colored according to edge weights.



Figure 5: Velocity relations graph built using the expression  $w_{ij}^{\cos,v} = \frac{1}{2} (1 + \cos \phi) \cdot \frac{1}{2} (|\mathbf{v}_i| + |\mathbf{v}_j|)$  (Eq. 6) for edge weight, with only the edges exceeding the weight threshold shown: at edge weight 0.05 (a) and 0.15 (b) from the maximum edge weight value.



Figure 6: Velocity relations graph built using the expression  $w_{ij}^{\cos,v} = \frac{1}{2} (1 + \cos \phi)$  (Eq. 5) for edge weight, with only the edges exceeding the weight threshold shown: at edge weight 0.05 (a) and 0.60 (b) from the maximum edge weight value.

Figure 7: Velocity relations graph built using the expression  $w_{ij}^{dot} = v_i \ v_j$ , (Eq. 3) for edge weight, with only the edges exceeding the weight threshold shown: at edge weight 0.05 (a) and 0.15 (b) from the maximum edge weight value.

Figure 9: Communities detected from the graph with the edge weight de ned as  $w_{ij}^{\cos;v} = \frac{1}{2}(1 + \cos) = \frac{1}{2}(jv_ij + jv_jj)$  (Eq. 6), shown as green outlines, superimposed on the velocity elds at di erent speed thresholds at t = 0:7 s (cf. Fig. 2).

## 5 CONCLUSIONS

We have presented an approach for identification of structures in the granular velocity fields consisting of groups of particles moving in a coordinated matter, based on the graph community detection algorithm. For this purpose, a graph is built with edge weights calculated based on velocity relationships between the adjacent particle pairs. The choice of a particular expression for calculation of the graph edge weights has a considerable influence to the resulting graph structure and consequently influences the results of the community detection. In the presented case, an expression for the graph edge weight based on the cosine of the angle between the velocity vectors of the respective particle pairs, taking into account their velocities, turned out to be best suited for our purposes. In many cases, it enables to identify the velocity field structures with high precision.

# 6 ACKNOWLEDGMENTS

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# 7 CONFLICTS OF INTEREST

This work does not have any conflicts of interest.

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